Journal of Applied Mathematics & Bioinformatics, vol. x, no. x, 2011, x-x ISSN: 1792-6602(print), 1792-6939(online) ©International Scientific Press, 2011

A try to prove Riemann's Hypothesis

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Abstract

Riemann's conjecture (1859) states that:

The real part of every non trivial zero of Riemann's zeta function is 1/2.

The main contribution of this paper is to use the classical approach to this conjecture whose the key idea is to provide a self-adjoint operator whose real eigenvalues are the imaginary part of the non trivial zeros of Riemann's zeta function and whose existence, according to Hilbert and Pólya, proves Riemann's hypothesis.

Mathematics Subject Classification : 11A41; 11M06; 11M26 Keywords: Euler; Riemann; Hilbert; Polya; conjecture

1 Introduction

In his book [1] of 1748, Leonhard Euler (1707-1783) proved what is now named *the Euler product formula*. This product is the result of the infinite sum:

$$\sum_{n=1}^{\infty} \frac{1}{n^s} = \prod_{p \in \mathbb{P}} (1 - 1/p^s)^{-1} \quad \text{for any integer variable } s > 1$$

Article Info: Received : January 10, 2011. Published online : July 17, 2011

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where \mathbb{P} is the infinite set of primes.

In his article [2] of 1859, Riemann (1826-1866) extended the Euler definition to the complex variable s of the zeta function:

$$\zeta(s) = \prod_{p \in \mathbb{P}} (1 - 1/p^s)^{-1}$$
 for any complex variable $s \neq 1$

It is known that the trivial zeros of the function are the infinite set:

 $\{s_1\} = -2m$ for all integers m > 0

Riemann's hypothesis can be seen as stating that:

Probably, the infinite set of the non trivial zeros $\{s_2\}$ of $\zeta(s)$ can be written:

 $\{s_2\} = \frac{1}{2} + it_n$ where t_n is real.

This conjecture is the first point of the eighth unresolved problem (among 23) that Hilbert listed in 1900 [3] as well as the second unresolved problem listed in 2000 by The Clay Mathematics Institute [4].

2 Preliminary notes

2.1 Hilbert-Pólya statement

Circa 1914, Hilbert et Pólya [5], independently from each other, have orally stated that Riemann's hypothesis would be proved if it could be shown that the imaginary parts t_n of the non trivial zeros of the symmetrical xi function $\xi(s)$ derived from $\zeta(s)$, corresponded to the real eigenvalues of an unbounded self-adjoint operator (here named \hat{H}_{ξ}) for which we could write:

$$H_{\xi}\psi_k = E_k\psi_k \tag{1}$$

which is an equation of quantum physics where E is for energy.

So, the first and unique purely mathematical clue that we have is that this operator should be a square matrix of infinite dimension with real eigenvalues. This means that it could be written:

$$\hat{H}_{\xi} = \begin{pmatrix} \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \\ \cdots & t_{n-1} & 0 & 0 & 0 & \cdots \\ \cdots & 0 & t_n & 0 & 0 & \cdots \\ \cdots & 0 & 0 & t_{n+1} & 0 & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix} \quad \text{all } t_i \text{ being real}$$

3 Proof of Riemann's Hypothesis

The proof will be established in two steps:

The first one establishes only a conditional proof.

The second one establishes the unconditional proof.

3.1 Conditional proof of Riemann Hypothesis

Proof. By definition, a complex number s is written:

$$s = x + iy$$
 where x and y are real and $i = \sqrt{-1}$

By changing the conventional basis of coordinates (x, y) of the complex plane into the new one $(x' = \frac{1}{2} - x, y' = y)$, these complex numbers can be written:

s'=x'+iy' in the new basis or: $s'=(\tfrac{1}{2}-x)+iy \text{ using the change in coordinates}.$

Condition. We suppose that the \hat{H}_{ξ} operator exists and that it contains the infinitely many real eigenvalues t_n coming from the non trivial zeros s_2 of $\zeta(s)$.

Hypothesis. We then suppose that these non trivial zeros lie anywhere in the complex plane with the two exceptions that they cannot lie on the real axis x or x' (reserved for trivial zeros s_1), which gives:

$$y \neq 0$$
 and $y' \neq 0$

nor on the conventional critical line $x = \frac{1}{2}$ that becomes the new imaginary axis y', which gives:

$$x \neq \frac{1}{2}$$
 and $x' \neq 0$

Then, each non trivial zero s_2 of $\zeta(s)$ could be written:

$$s_2 = x'_2 + iy'_2$$
 with $x'_2 \neq 0$ and $y'_2 \neq 0$
or, using the change in coordinates:
 $s_2 = (\frac{1}{2} - x_2) + iy_2$ with $x_2 \neq \frac{1}{2}$ and $y_2 \neq 0$

But using the fact that $-x_2 = i^2 x_2$, they can be written:

$$s_{2} = (\frac{1}{2} + i^{2}x_{2}) + iy_{2} = \frac{1}{2} + i(y_{2} + ix_{2}) \text{ with } x_{2} \neq \frac{1}{2} \text{ and } y_{2} \neq 0$$

or:
$$s_{2} = \frac{1}{2} + it_{2} \text{ with } t_{2} = y_{2} + ix_{2}, x_{2} \neq \frac{1}{2} \text{ and } y_{2} \neq 0$$

and we get the result that the non trivial zeros can exist only when $x_2 = 0$ to have t_2 real, a result that is false because it is known that all non trivial zeros lie in the critical strip 0 < x = Re(s) < 1 that excludes x = 0.

We thus get the contradiction to our hypothesis that the quantity $t_2 = y_2 + ix_2$ is unable, when the critical line $x_2 = \frac{1}{2}$ is excluded, to provide any real value $t_n \neq 0$ to the operator \hat{H}_{ξ} . Therefore, as it is known on one hand that infinitely many non trivial zeros do exist for $\zeta(s)$ and on the other hand that they cannot lie on the x axis nor out of the critical line $x = \frac{1}{2}$, it proves that they can only lie on this critical line. This, in turn, proves Riemann's hypothesis, conditionally to the existence of the \hat{H}_{ξ} operator.

3.2 Unconditional proof of Riemann's Hypothesis

Proof. Now, as Riemann's Hypothesis is conditionally proved, the \hat{H}_{ξ} operator conditionally exists and contains the infinitely many real values t_n of the non trivial zeros $\{s_2\}$ of $\zeta(s)$.

To prove that the \hat{H}_{ξ} operator *do exists*, we will consider the new and larger operator \hat{H}_{ζ} built with the zeros of both sets of zeros $\{s_1\}$ and $\{s_2\}$ of $\zeta(s)$ as eigenvalues, an operator that also contains the real values t_n but not as eigenvalues:

$$\hat{H}_{\zeta} = \begin{pmatrix} \dots & \dots \\ \dots & -6 & 0 & 0 & 0 & 0 & 0 & \dots \\ \dots & 0 & -4 & 0 & 0 & 0 & 0 & \dots \\ \dots & 0 & 0 & -2 & 0 & 0 & 0 & \dots \\ \dots & 0 & 0 & 0 & \frac{1}{2} + it_1 & 0 & 0 & \dots \\ \dots & 0 & 0 & 0 & 0 & \frac{1}{2} + it_2 & 0 & \dots \\ \dots & 0 & 0 & 0 & 0 & 0 & \frac{1}{2} + it_3 & \dots \\ \dots & \dots & \dots & \dots & \dots & \dots & \dots \end{pmatrix}$$

As this new operator contains the real values t_n , it enables us, at any time but if it exists, to rebuild the operator \hat{H}_{ξ} of Hilbert and Pólya. To simplify the writing, we set:

$$\begin{pmatrix} \dots & \dots & \dots & \dots \\ \dots & -6 & 0 & 0 \\ \dots & 0 & -4 & 0 \\ \dots & 0 & 0 & -2 \end{pmatrix} = (-2m)$$

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and:

$$\begin{pmatrix} \frac{1}{2} + it_1 & 0 & 0 & \dots \\ 0 & \frac{1}{2} + it_2 & 0 & \dots \\ 0 & 0 & \frac{1}{2} + it_3 & \dots \\ \dots & \dots & \dots & \dots \end{pmatrix} = \left(\frac{1}{2} + it_n\right)$$

so that \hat{H}_{ζ} can be written:

$$\hat{H}_{\zeta} = \begin{pmatrix} (-2m) & (0) \\ (0) & \left(\frac{1}{2} + it_n\right) \end{pmatrix}$$

But the matrices (-2m) and $(\frac{1}{2} + it_n)$ representing the sets of zeros $\{s_1\}$ and $\{s_2\}$ can symbolically be replaced by their parametric form:

$$-2m$$
, *m* being a positive integer parameter
 $\frac{1}{2} + it_n$, t_n being a real parameter

The sets $\{s_1\}$ and $\{s_2\}$ can then be considered as the two infinite sets of roots of the polynomial of complex variable s:

$$P(s) = (s - s_1)(s - s_2) = s^2 - (s_1 + s_2)s + s_1s_2$$

$$P(s, m, t) = s^2 - (-2m + \frac{1}{2} + it_n)s - 2m(\frac{1}{2} + it_n)$$

$$P(s, m, t) = s^2 + (2m - (\frac{1}{2} + it_n))s - 2m(\frac{1}{2} + it_n)$$

which can be written either:

$$P(s,m,t) = \begin{pmatrix} s^2 & s & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & (2m - (\frac{1}{2} + it_n)) & 0 \\ 0 & 0 & -2m(\frac{1}{2} + it_n) \end{pmatrix} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$
(2)

or:

$$P(s,m,t) = \begin{pmatrix} 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & (2m - (\frac{1}{2} + it_n)) & 0 \\ 0 & 0 & -2m(\frac{1}{2} + it_n) \end{pmatrix} \begin{pmatrix} s^2 \\ s \\ 1 \end{pmatrix}$$
(3)

Equation (2) gives, by multiplying the first two matrices:

$$P(s,m,t) = \left(s^2 \quad (2m - (\frac{1}{2} + it_n))s \quad -2m(\frac{1}{2} + it_n)\right) \begin{pmatrix} 1\\ 1\\ 1 \end{pmatrix} = E_k \psi_{E_k} \qquad (4)$$

when we set:

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$$E_{k} = \begin{pmatrix} s^{2} & (2m - (\frac{1}{2} + it_{n}))s & -2m(\frac{1}{2} + it_{n}) \end{pmatrix}$$

and:
$$\psi_{E_{k}} = \begin{pmatrix} 1\\ 1\\ 1 \end{pmatrix}$$

Now, as from (4) we also have:

$$P(s,m,t) = \left(1 \quad (2m - (\frac{1}{2} + it_n)) \quad -2m(\frac{1}{2} + it_n)\right) \begin{pmatrix} s^2 \\ s \\ 1 \end{pmatrix} = H_k \psi_{H_k} \qquad (5)$$

when we set:

$$H_{k} = \begin{pmatrix} 1 & (2m - (\frac{1}{2} + it_{n})) & -2m(\frac{1}{2} + it_{n}) \end{pmatrix}$$

and:
$$\psi_{H_{k}} = \begin{pmatrix} s^{2} \\ s \\ 1 \end{pmatrix} = \hat{R} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = \hat{R}\psi_{E_{k}}$$

where \hat{R} is the 3-dimensional rotation matrix from the orthogonal basis ψ_{H_k} used to describe H_k to the orthogonal basis ψ_{E_k} used to describe E_k , and we have:

$$H_k \psi_{H_k} = H_k \hat{R} \psi_{E_k} = E_k \psi_{E_k}$$

Then, setting $\hat{H} = H_k \hat{R}$, we get:

$$\hat{H}\psi_{E_k} = E_k \psi_{E_k} \tag{6}$$

which is identical to equation (1). And the \hat{H}_{ξ} operator looked for by Hilbert and Pólya can be built with the real values t_n of the existing operator:

$$\hat{H} = H_k \hat{R} = \left(1 \left(2m - \left(\frac{1}{2} + it_n\right) \right) - 2m\left(\frac{1}{2} + it_n\right) \right) \hat{R}$$

As we can rebuild the self-adjoint operator \hat{H}_{ξ} linked to $\zeta(s)$ via the function P(s, m, t) and the existing operator \hat{H} , this self-adjoint operator \hat{H}_{ξ} do exists and as we have proved earlier that Riemann hypothesis is true conditionally to the existence of the \hat{H}_{ξ} operator, Riemann hypothesis is unconditionally proved.

Remark. As H_k can also be written:

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$$H_k = \begin{pmatrix} 1 & 1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & (2m - (\frac{1}{2} + it_n)) & 0 \\ 0 & 0 & -2m(\frac{1}{2} + it_n) \end{pmatrix} = \begin{pmatrix} 1 & 1 & 1 \end{pmatrix} \hat{A}$$

when we set:

$$\hat{A} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & (2m - (\frac{1}{2} + it_n)) & 0 \\ 0 & 0 & -2m(\frac{1}{2} + it_n) \end{pmatrix}$$
(7)

we get from (3) and (2) that:

$$\begin{pmatrix} 1 & 1 & 1 \end{pmatrix} \hat{A} \begin{pmatrix} s^2 \\ s \\ 1 \end{pmatrix} = P(s, m, t) = \begin{pmatrix} s^2 & s & 1 \end{pmatrix} \hat{A} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$$
(8)

But, for any s = x + iy, we have:

$$P(s,m,t) = s^{2} + \left(2m - \left(\frac{1}{2} + it_{n}\right)\right)s - 2m\left(\frac{1}{2} + it_{n}\right)$$

$$= (x + iy)^{2} + \left(2m - \left(\frac{1}{2} + it_{n}\right)\right)(x + iy) - 2m\left(\frac{1}{2} + it_{n}\right)$$

$$= \left(x^{2} - y^{2} + \left(2m - \frac{1}{2}\right)x + yt_{n}\right) + i\left(-xt_{n} + y\left(2m - \frac{1}{2}\right)\right) - m - 2mit_{n}$$

$$= \left(x^{2} - y^{2} + \left(2m - \frac{1}{2}\right)x + yt_{n} - m\right) + i\left(-xt_{n} + y\left(2m - \frac{1}{2}\right) - 2mt_{n}\right)$$

and P(s, m, t) will be real only when:

 $-xt_n + y(2m - \frac{1}{2}) - 2mt_n = 0$ and so, for all of the infinitely many lines in the complex plane such that: $y = \frac{t_n}{2m - \frac{1}{2}}x + \frac{2mt_n}{2m - \frac{1}{2}}$ where *m* and *t_n* are the discrete values defined earlier

Then, for all the points of all these lines, we have:

$$P(s,m,t)_{lines} = \left(x^2 - y^2 + (2m - \frac{1}{2})x + yt_n\right) = V(x,y), \text{ a real value}$$
(9)

and therefore for the mono-term matrix (V(x,y)) we have:

$$(V(x,y)) = (\overline{V(x,y)}) = (\overline{V(x,y)})^T$$
(10)

where $(\overline{V(x,y)})$ is the conjugate matrix of (V(x,y)) and $(\overline{V(x,y)})^T$ is the conjugate transpose of (V(x, y)). So, from (8), (9) and (10), we can write:

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$$P(s,m,t)_{lines} = V(x,y) = \begin{pmatrix} 1 & 1 & 1 \end{pmatrix} \hat{A} \begin{pmatrix} s^2 \\ s \\ 1 \end{pmatrix} = \left(\overline{\begin{pmatrix} s^2 & s & 1 \end{pmatrix} \hat{A} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}} \right)^T$$

which proves that the operator \hat{A} that also provides the real t_n to \hat{H}_{ξ} , verifies the equation of the observables in quantum physics, which is generally written:

$$\langle \psi_1 \mid \hat{A} \mid \psi_2 \rangle = \left(\langle \psi_2 \mid \hat{A} \mid \psi_1 \rangle \right)^T$$

where \hat{A} is a self-adjoint operator associated to a physical quantity A, $\langle \mathbf{x} |$ and $|\mathbf{x}\rangle$ are the bra and ket operators and ψ_1 and ψ_2 are the states of the physical quantity A before and after the measuring of A.

ACKNOWLEDGEMENTS. This work is dedicated to my family. As I am a hobbyist, I wish to express my gratitude towards the Editors of this journal as well as towards the team of Reviewers for having welcomed, reviewed and accepted my article.

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